



Raman Spectroscopy of the $\text{Ni}_5(\text{TeO}_3)_4\text{Cl}_2$ Single Crystal

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The two-dimensional quantum spin system $\text{Ni}_5(\text{TeO}_3)_4\text{Cl}_2$ has a clearly pronounced 2D spin structure with the developed Ni^{2+} triangular structure that can give rise to the appearance of the exotic magnetic states. The antiferromagnetic transition was previously revealed in the single crystal at temperature $T_N = 28.5$ K [1]. Raman spectroscopy is the informative non-destructive method that can probe the interplay of lattice, charge, and spin quantum excitations. In our previous work [2], the main attention was paid to the analysis of the phonon spectrum. In the present work, we focus on the analysis of the revealed Raman peculiarities of the $\text{Ni}_5(\text{TeO}_3)_4\text{Cl}_2$ single crystal.

The crystal $\text{Ni}_5(\text{TeO}_3)_4\text{Cl}_2$ is monoclinic with space group $C2/c$ (C_{2h}^6) and $Z = 4$. The crystal unit cell has a dimension: $a=19.5674\text{\AA}$, $b=5.2457\text{\AA}$, $c=16.3084\text{\AA}$, and $\beta=125.29^\circ$ [3].

Raman spectra of the $\text{Ni}_5(\text{TeO}_3)_4\text{Cl}_2$ single crystal presented in the present work were taken in the frequency region of 3–850 cm^{-1} in the temperature range of 2–300 K. A_g and B_g modes are active in Raman tensor components: $A_g - XX, YY, ZZ, XY, YX, B_g - XZ, ZX, YZ, ZY$.

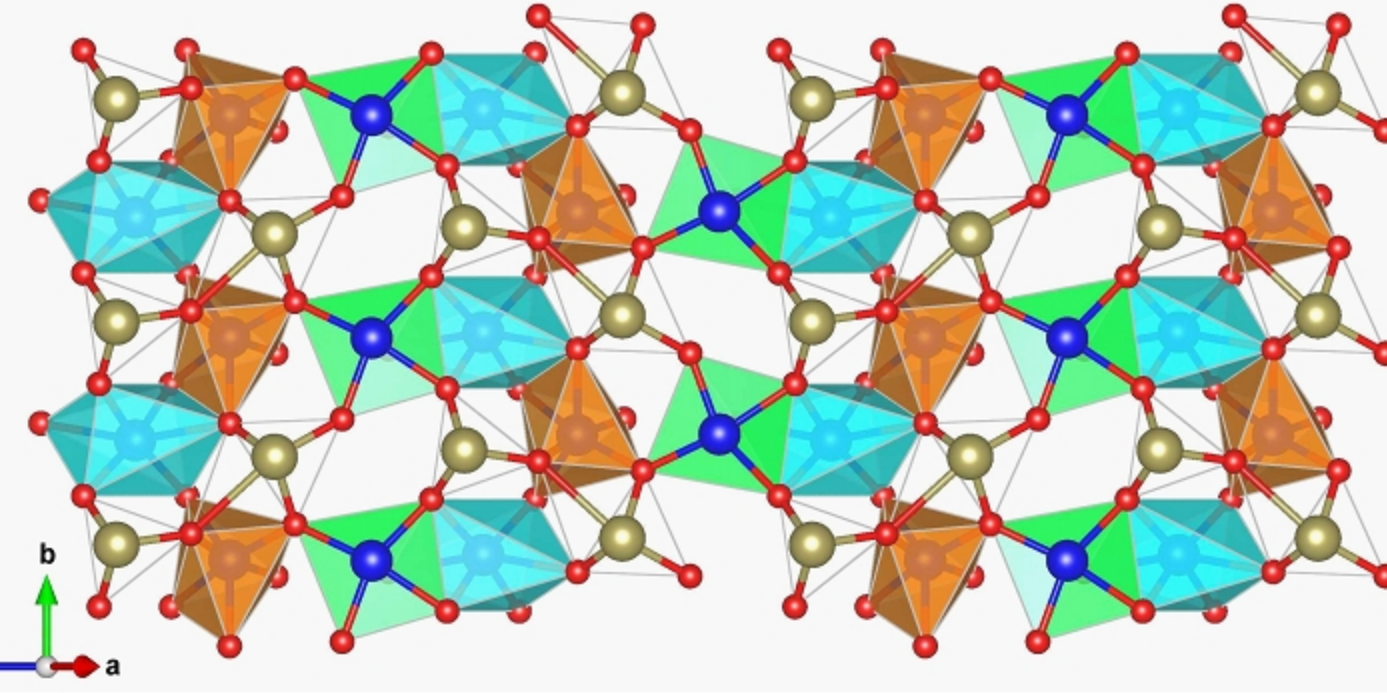


Fig. 1. Ni octahedral shaded layer [1]. Ti atoms are shown in mustard color.

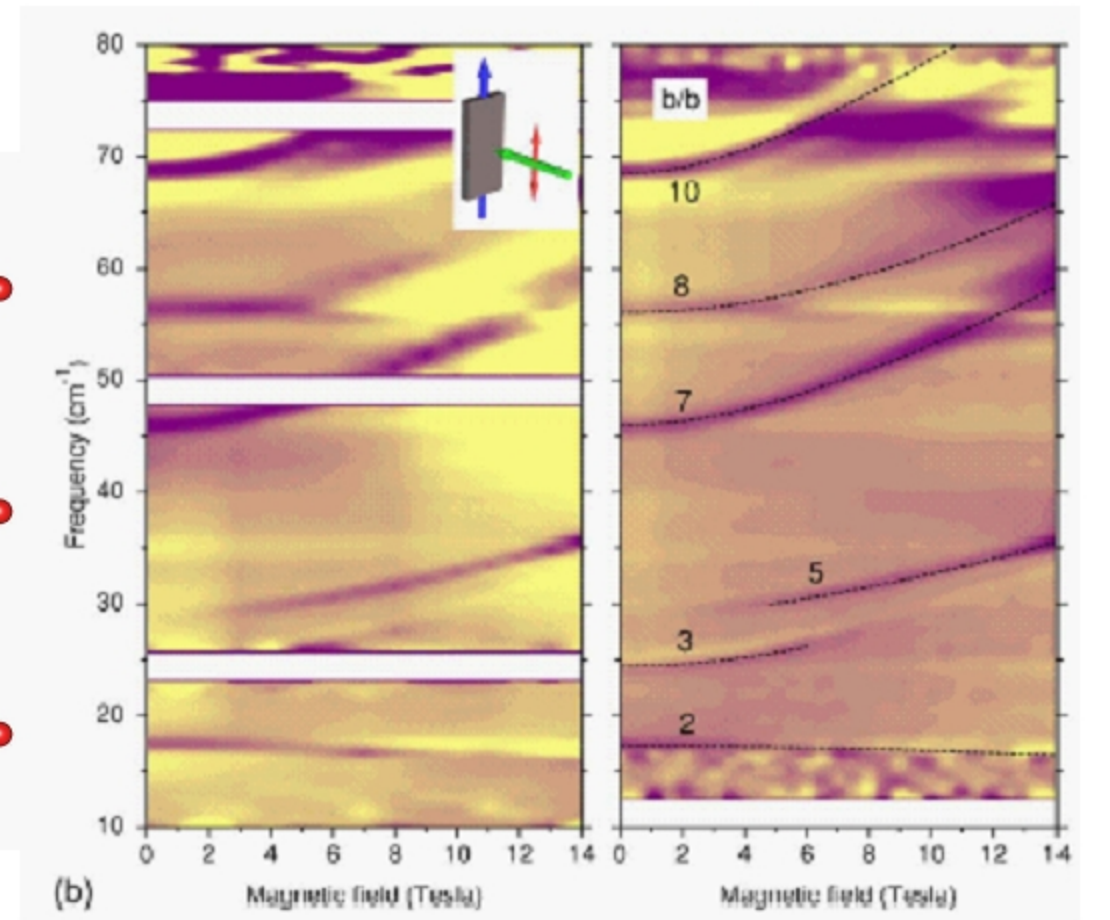


Fig. 2. Spin resonance in the ordered magnetic state of $\text{Ni}_5(\text{TeO}_3)_4\text{Cl}_2$ [4].

Previous Raman work [2]

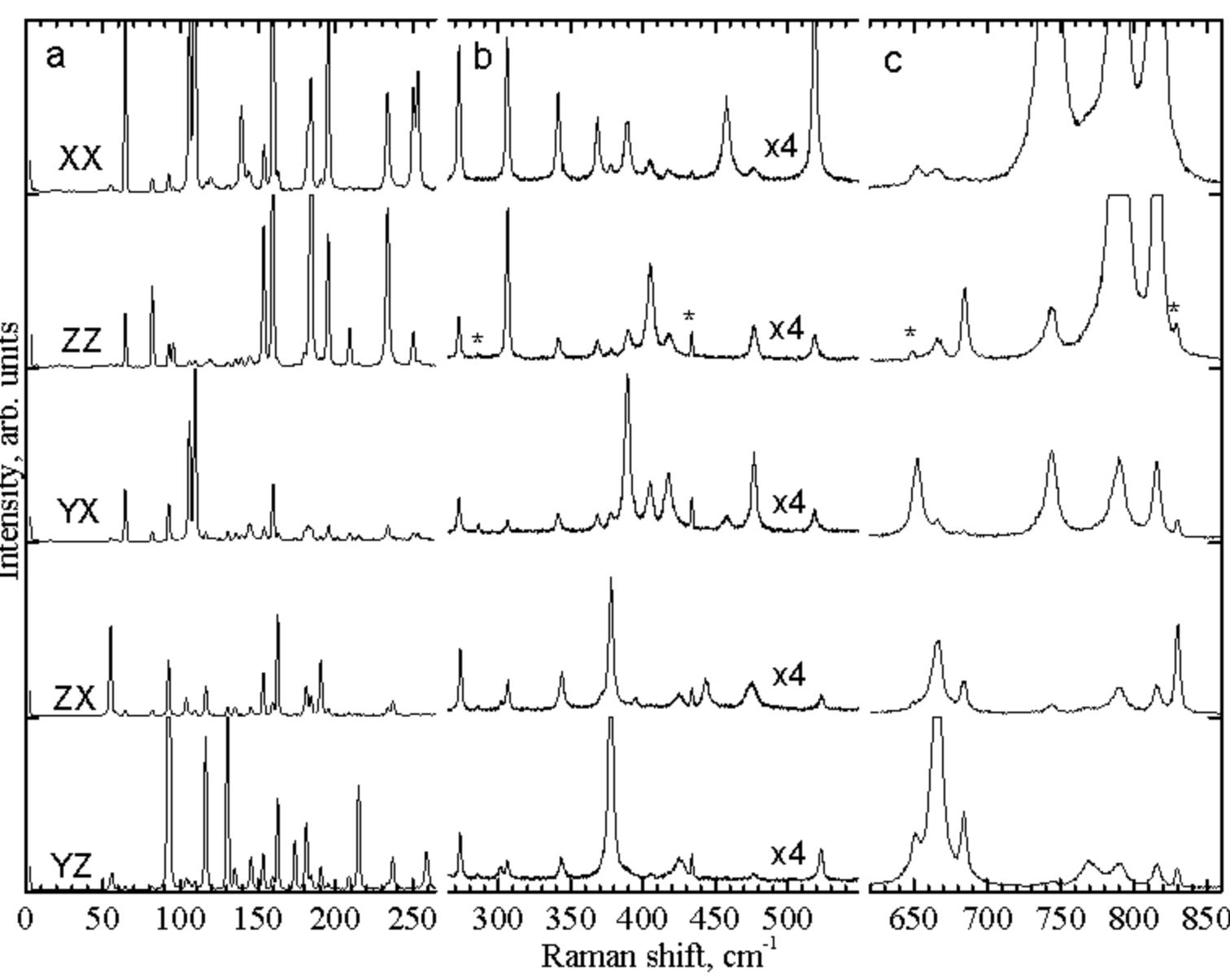


Fig. 3, 4. Raman spectra at 15 K (25 K) [2].

Frequency (cm^{-1})	Symmetry	Experimental scattering geometry	Frequency (cm^{-1})	Symmetry	Experimental scattering geometry	Free TeO_3 unit [10]
55.1	B_g	ZZ	253.4	A_g	XX	$\nu_4(E)$ 326
64.5	A_g	XX	259.0	B_g	YZ	
81.8	A_g	ZZ	272.8	A_g	XX	$\nu_2(A_1)$ 364
92.5	B_g	YZ	273.8	B_g	ZZ, YZ	
95.4	A_g	ZZ	301.8	B_g	ZZ	$\nu_3(E)$ 703
103.8	B_g	ZZ	306.3	A_g	ZZ, XX	
105.8	A_g	ZZ, XX	341.5	A_g	XX	$\nu_1(A_1)$ 758
109.4	A_g	XX, YX	343.8	B_g	YZ	
116.5	B_g	YZ	368.5	A_g	XX	
119.6	A_g	XX	377.9	B_g	YZ	
130.5	B_g	YZ	389.3	A_g	XX, YX	
134.8	B_g	YZ	395.0	B_g	ZZ	
139.2	A_g	XX	405.0	A_g	ZZ	
144.5	A_g	XX, YX	417.9	A_g	YX	
145.6	B_g	YZ	425.0	B_g	ZZ, YZ	
153.5	B_g	ZZ, YZ	443.7	B_g	ZZ	
153.8	A_g	ZZ	457.9	A_g	XX	
159.7	A_g	ZZ, XX	474.6	B_g	ZZ	
162.9	B_g	YZ, YZ	476.8	A_g	YX	
174.0	B_g	ZZ, YZ	518.8	A_g	XX	
181.4	B_g	YZ	523.2	B_g	YZ	
182.5	A_g	XX	651.1	B_g	YZ	
184.4	A_g	ZZ	652.3	A_g	YX	
190.7	B_g	ZZ	666.0	B_g	YZ	
195.4	A_g	ZZ, XX	684.2	B_g	YZ	
209.0	B_g	YZ	684.9	A_g	ZZ	
209.4	A_g	ZZ	743.9	A_g	XX	
215.1	B_g	YZ	770.1	B_g	YZ	
233.7	A_g	ZZ	789.9	A_g	ZZ, XX	
237.1	B_g	YZ	815.6	A_g	ZZ, XX	
250.3	A_g	XX	829.9	B_g	ZZ	

$$\Gamma_{\text{vib}} = 34A_g + 35B_g + 34A_u + 35B_u$$

$34A_g + 35B_g$ modes are Raman active.

In different scattering geometries we observed 62 phonon lines (Table 1) of the 69 predicted ones by theory-group analysis for the monoclinic C_{2h}^6 lattice symmetry [2]. Three additional weak lines with frequencies 22.3, 29.4, and 49.0 cm^{-1} vanish at around T_N with temperature rising [2].

References:

- [1] R. Becker, M. Prester, H. Berger, M. Johansson, D. Drobac, and I. Zivkovic, Sol. State Sci. 9, 223 (2007).
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- [3] M. Johansson, K.W. Tornroos, P. Lemmens, and P. Millet, Chem. Mater. 15, 68 (2003).
- [4] Mihaly, T. Feher, B. Dora, B. Nafraadi, H. Berger, and L.Ferro, Phys. Rev. B 74, 174403 (2006).

Table 1. Symmetry and frequencies of the phonon modes in $\text{Ni}_5(\text{TeO}_3)_4\text{Cl}_2$ [2] at $T = 15$ K.

New Raman data

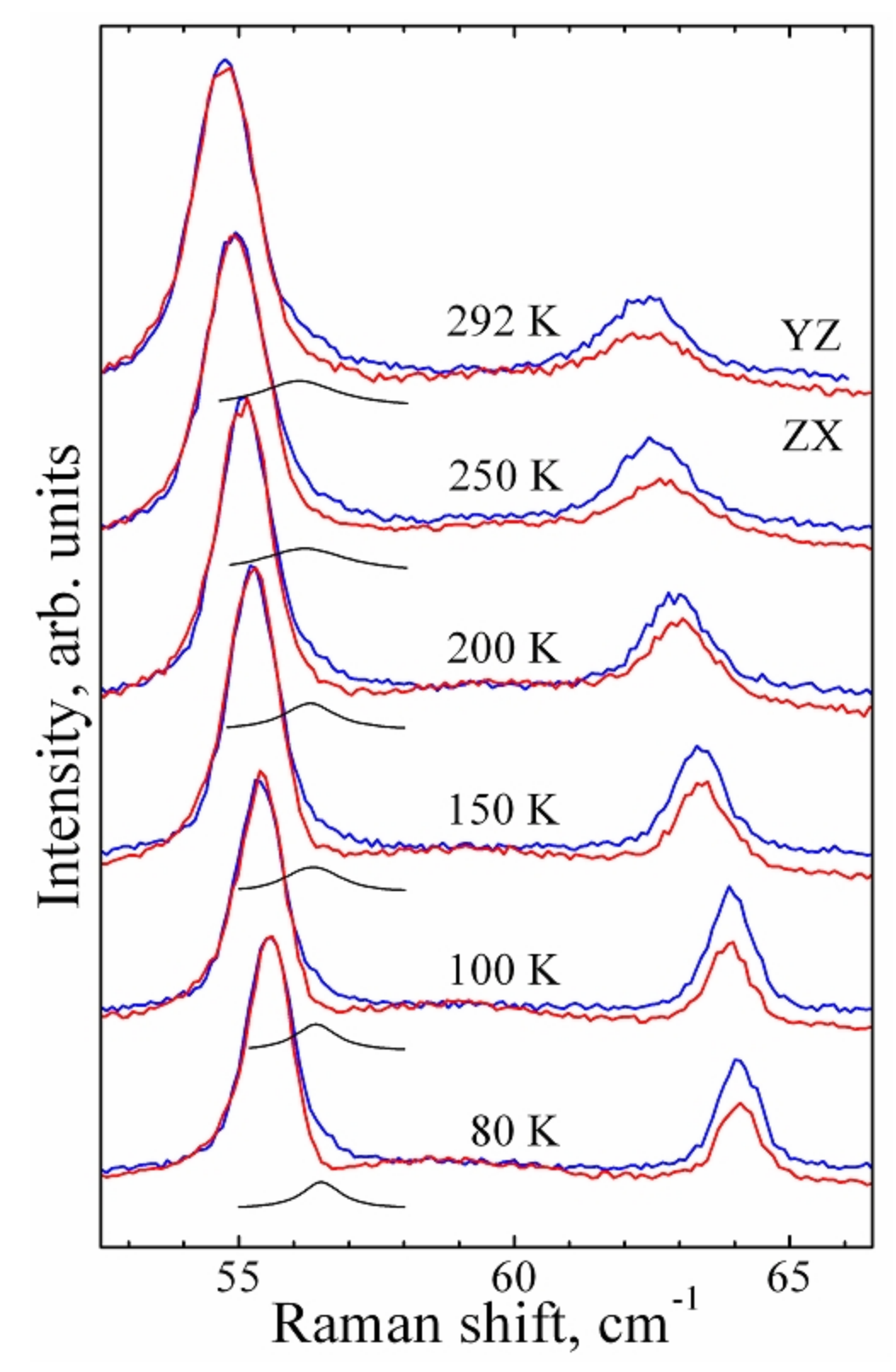
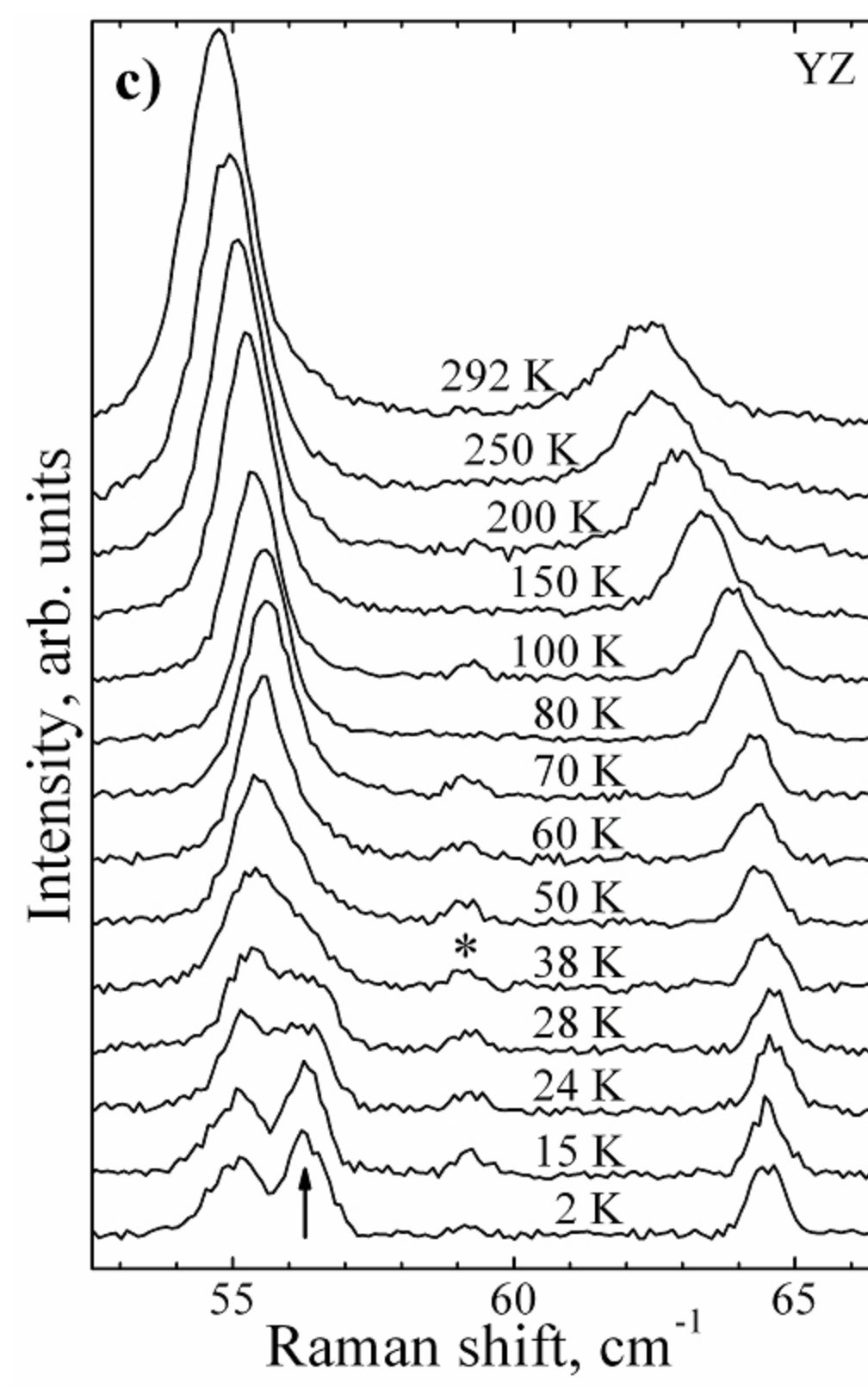
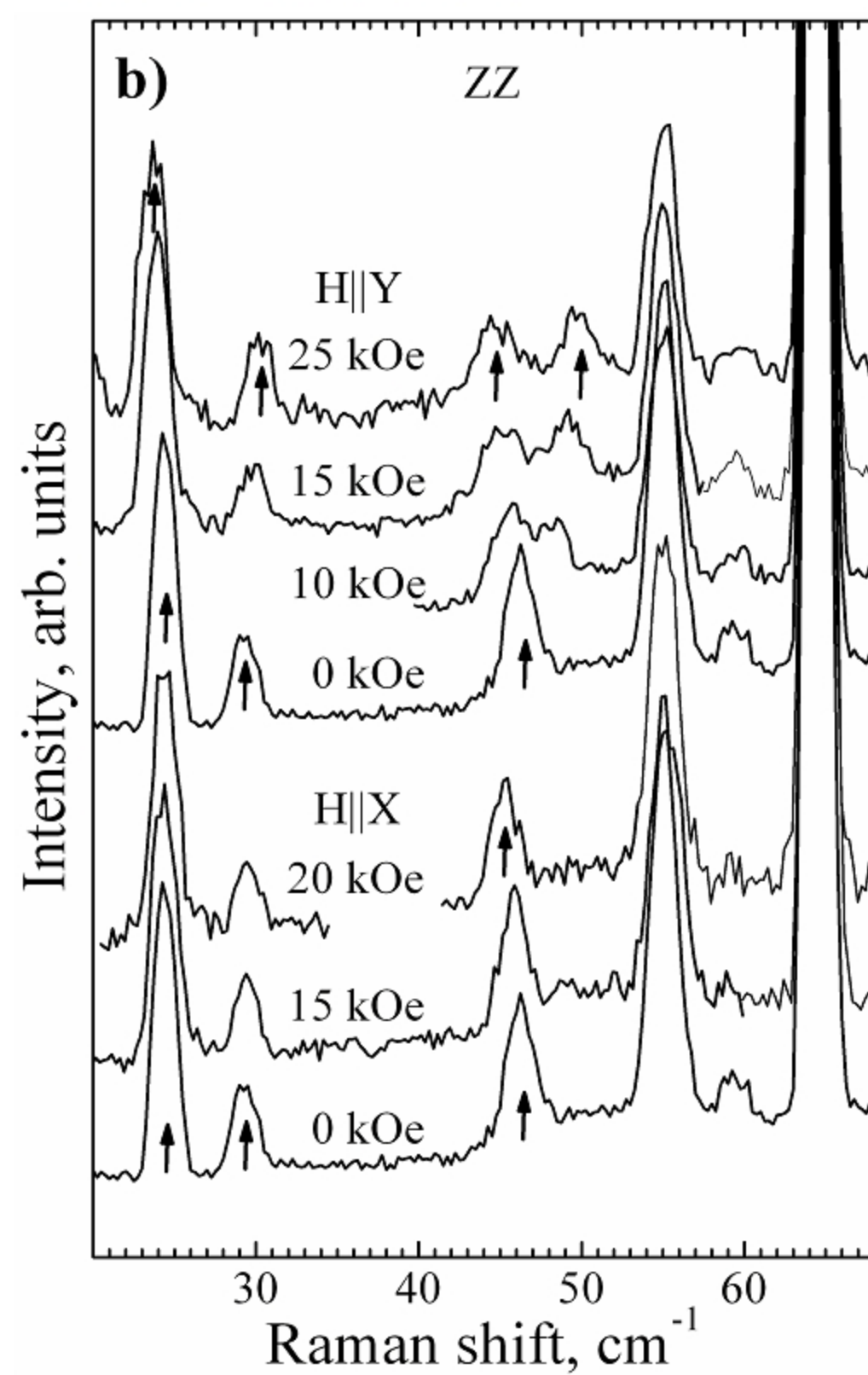
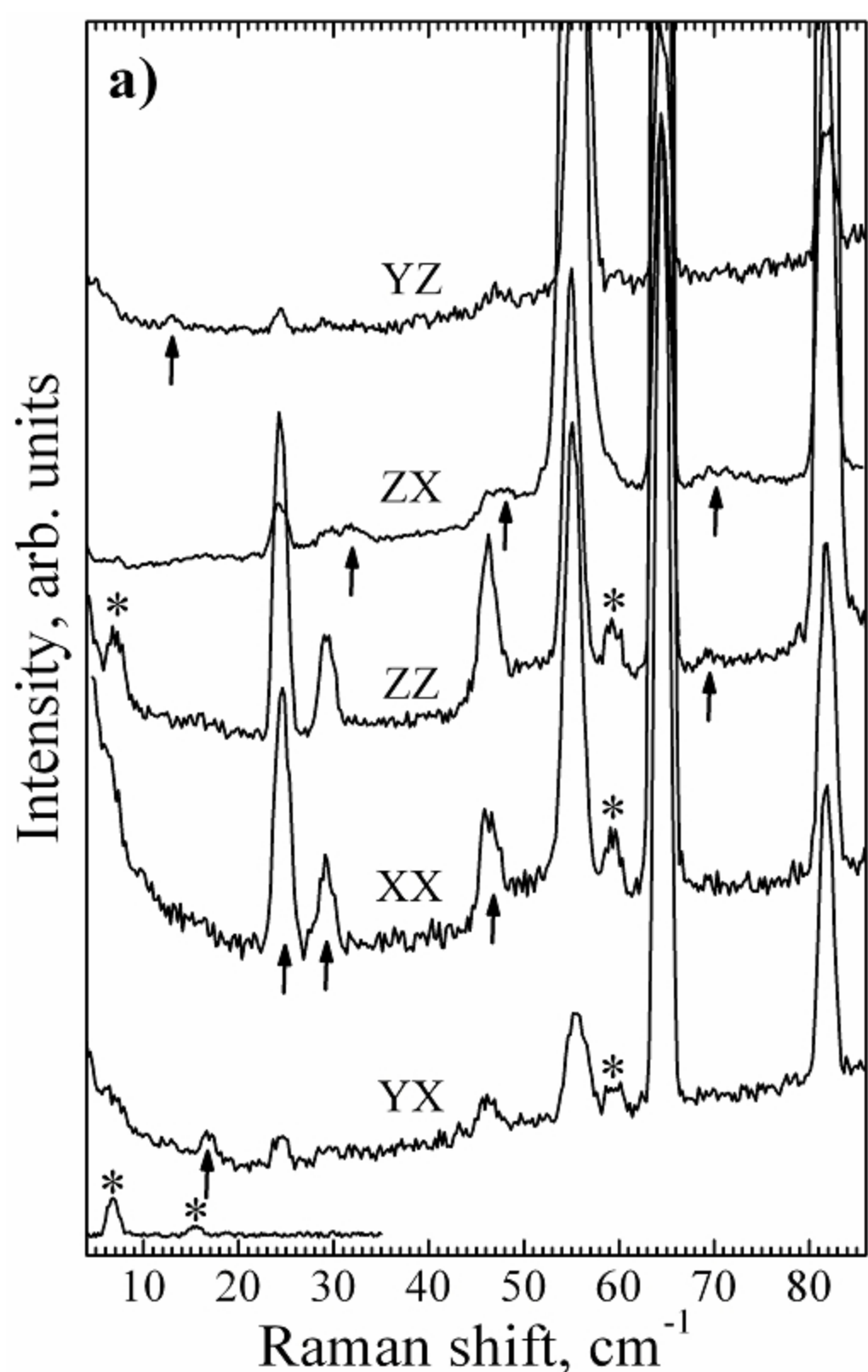


Fig. 5. a) Low-frequency polarized Raman spectra at $T = 2$ K. b) Dependence of Raman spectra on the applied magnetic fields (H) at $T = 2$ K. c) Temperature dependent Raman spectra. Spectral resolution: 1.8 cm^{-1} (a and b); 0.6 cm^{-1} (c). The arrows mark additional lines. The asterisks denote the plasma lines of He-Ne laser ($\lambda_{\text{exc}} = 632.8$ nm).

Fig. 6. Temperature dependence of B_g (YZ and ZX) Raman spectra. Spectral resolution: 0.6 cm^{-1} . $\lambda_{\text{exc}} = 632.8$ nm.

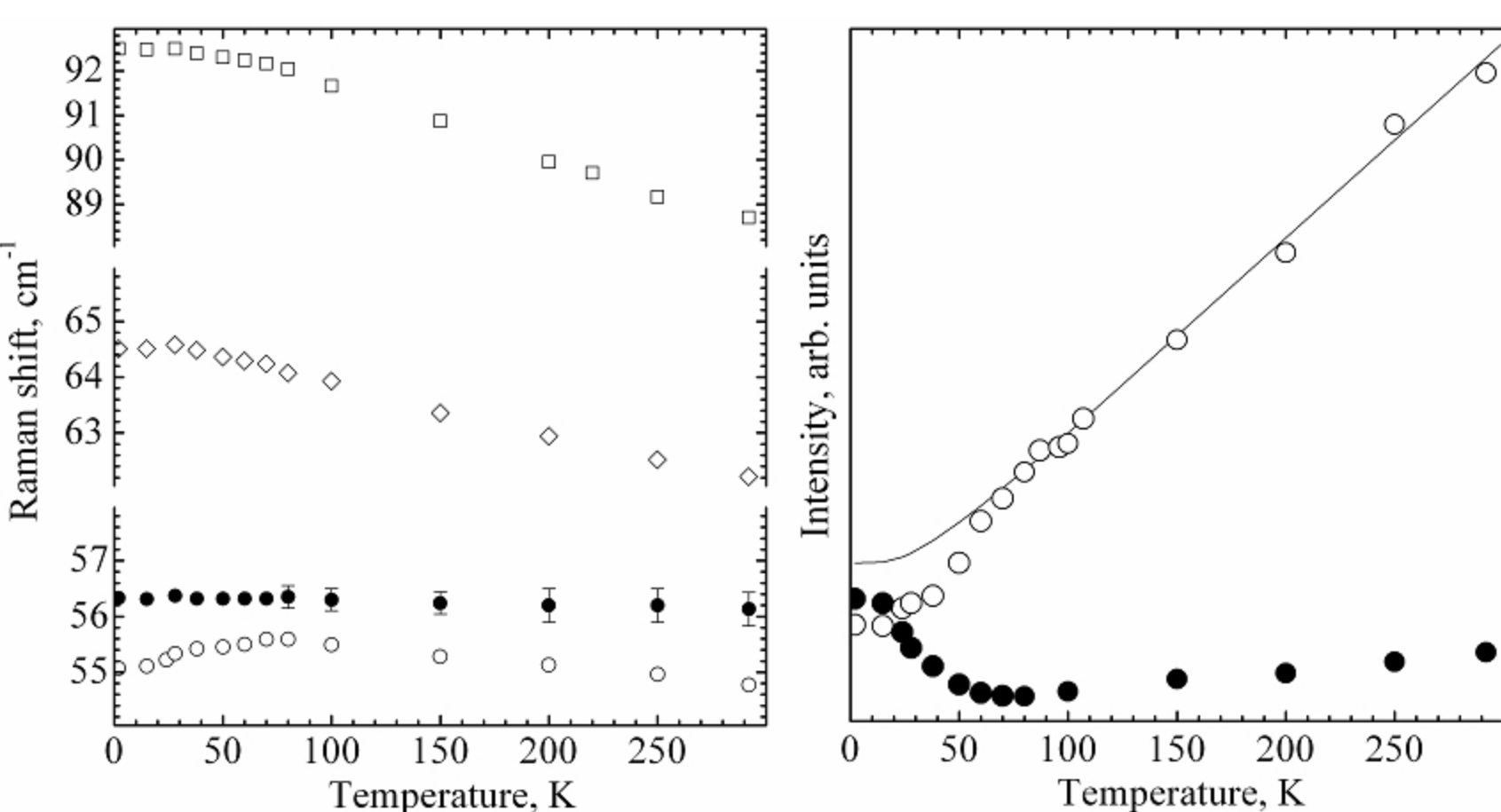


Fig. 7. Temperature dependence of the Raman position (left panel) and intensity (right panel) of the phonon modes: \circ - 55 cm^{-1} , \bullet - 56.3 cm^{-1} , \diamond - 64.5 cm^{-1} , \square - 92.5 cm^{-1} . The spectral resolution: 0.6 cm^{-1} .

Raman scattering data		FIR (2.5 K) [4]	
Previous work (15 K) [2]	Present work (2 K)	At $H = 0$	At high field
	13.1±0.5	13.5	
	16.8±0.5	17.3	
22.3±0.5	24.4±0.2	24.4	
			25.5
29.4±0.5	29.2±0.3	29	
	31.9±0.4	32.2	
49.0±1.0	46.2±0.2	46	
	48.2±0.5	46	
56.3±0.1?	56.3±0.1 phonon	56	
			63
	69.4±0.4	68.5	
	69.9±0.5	68.5	
two-magnon	50.0±1.0		

Conclusions

In addition to the phonon lines, weak peaks at 16.8, 24.4, 29.2, 46.2, 69.4 cm^{-1} (A_g modes) and 13.1, 31.9, 48.2, 69.9 cm^{-1} (B_g modes) were observed in the spectra at 2 K (Fig. 5a). The energy positions of additional peaks are in good accordance with spin-resonance data [4]. The temperature behavior of these lines is typical for the one-magnon scattering. The evolution of additional peaks has been investigated in applied magnetic field of 0 – 25 kOe in different geometries (Fig. 5b).

We have found that the 56.3 cm^{-1} line disappears in the spectra at temperatures well above T_N (Fig. 5c, 6). Besides, this line is not sensitive to the applied magnetic field. Usually such behavior is typical for a two-magnon scattering. However, the temperature dependence of its linewidth (see Fig. 5c) does not allow us to attribute this excitation to either one-magnon or two-magnon. We suppose that this line can be assigned to a phonon excitation.